

EXPERIMENT ON CRACK PROPAGATION IN A
VISCOELASTIC STRIP

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**CASE FILE
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INTRODUCTION AND SUMMARY

Crack propagation in viscoelastic material lies over two fields of solid mechanics, fracture mechanics and viscoelasticity. The classical principle of fracture mechanics "Griffith Theory" gives the condition for crack propagation and the application of viscoelasticity gives time-dependence of crack growth. This research treats crack propagation in a viscoelastic strip with an initial crack at the center.

The first part of this report is an abstract from the research of W. G. Knauss [1] and H. K. Mueller [2]. This part includes the energy equation, stress analysis and failure criteria for elastic material. These theories are expanded to viscoelastic material by using the Correspondence Principle.

The second part includes illustrations of this experiment.

The comparison between theory and experiment is in the third part of this report.

The failure criterion of a linearly viscoelastic material is discussed by W. G. Knauss and the crack propagation with constant velocity is discussed by H. K. Mueller. The purpose of this experiment is investigation of the possibility of determining crack velocity for short crack lengths where crack velocity is not constant upon the foundation of theories mentioned above.

1. THEORETICAL FOUNDATION

Let us consider a tip of a moving crack which is blunt. It may be reasonable to assume existence of a small region where the material is no longer continuous but forms a gathering of filaments.

A simple model of a crack tip is shown in Figure 1. Suppose the filamentized region is so small that this region has no effect on the entire stress strain field and these filaments give surface traction.

Let the length of this region be Δa and its width infinitesimally small. The amount of surface traction is equal to that of a continuously elastic medium which has no separated region Δa .

The crack travels a distance of Δa in the time Δt . A point P at position ξ on x-axis moves vertically to the point P' at the distance U_y . Figure 2 shows this process. The solid line indicates original position and the dotted line indicates the crack position after crack motion and this transition of the crack position creates a new open surface Δa .

The energy, ΔE_c , to create new surface Δa is equal to the work done by the vertical movement of new surface against the surface traction σ_y ,

$$-\Delta E_c = z \int_{\xi=a}^{a+\Delta a} \frac{1}{2} \sigma_y(\xi) U_y(\xi - \Delta a) d\xi \quad (1)$$

The surface traction at position U_y is zero and at center line is σ_y , the relation between surface traction and vertical displacement is supposed to be linear.

Linear elastic material in plane stress situation shows the following stress σ_y and displacement U_y^e in the vicinity of a crack tip.

$$\sigma_y(x, 0) = \frac{K_n}{\sqrt{\frac{x-a}{b}}} \sigma_o \quad \text{for } x \geq a \quad (2)$$

$$\sigma_y(x, 0) = 0 \quad \text{for } x < a$$

$$U_y^e(x, 0) = 4K_n \frac{\sigma_o}{E} \sqrt{b(a-x)} \quad \text{for } x \leq a \quad (3)$$

$$U_y^e(x, 0) = 0 \quad \text{for } x > a$$

where a : half length of the crack

b : width of the crack

K_n : stress intensity factor. Fig. 3

σ_o : stress at the infinite distance.

By substituting equation (2) and (3) into (1), one can get

$$\begin{aligned} & \int_{x=a}^{a+\Delta a} \sigma_y(x) U_y^e(x-\Delta a) dx \\ &= 4 \frac{\sigma_o^2 K_n^2}{E} b \int_{x=a}^{a+\Delta a} \sqrt{\frac{(a+\Delta a)-x}{x-a}} dx. \end{aligned} \quad (4)$$

Let us change variables $(x-a)/\Delta a = \cos^2 \theta$,

$$\int_{x=a}^{a+\Delta a} \sqrt{\frac{(a+\Delta a)-x}{x-a}} = -\int_{\frac{\pi}{2}}^0 \sqrt{\frac{1}{\cos^2 \theta}} -1 \Delta a \sin \theta \cos \theta d\theta$$

$$= -\frac{\pi \Delta a}{2}.$$

Thus,

$$\Delta E_c = 2\pi \frac{\sigma_o^2}{E} K_n^2 b \Delta a \quad (5)$$

Equation (5) leads the critical stress for crack propagation to the following expression

$$2S \leq \frac{\Delta E_c}{\Delta a} = \frac{\sigma_o}{E} K_n^2 b \pi \quad (6)$$

$$\sigma_{o \text{ crit.}} = \sqrt{\frac{S E}{\pi b K_n^2}}$$

S is the surface energy which is necessary for creating a new unit surface.

Equation (6) corresponds to the Griffith criterion for an infinitely large plate with crack under uniaxial tension when b tends to be infinitely large. That is,

$$\sigma_{o \text{ crit.}} = \sqrt{\frac{2 S E}{\Delta a}}.$$

The displacement U_y for viscoelastic materials is given by equation (7) by virtue of the Correspondence Principle [3].

$$U_y(x, o, t) = D_{cr}(t) E U_y^e(x, o) \quad (7)$$

where D_{cr} is the creep function.

The crack half length a grows as time goes on, so one can say U_y^e is a function of time. Now one can get displacement U_y in the case of moving crack.

$$U_y(x, 0, t) = D_{cr}(t) U_y^e(x, 0) + \int_0^t \frac{\partial}{\partial a} (U_y^e) \cdot \frac{\partial a(\tau)}{\partial \tau} D_{cr}(t-\tau) d\tau \quad (8)$$

Supposing the crack length is zero at $t=0$ and

$$\frac{\partial U_y^e}{\partial a} = 2 \frac{\sigma_0}{E} \sqrt{\frac{b}{a-x}} K_n(a/b),$$

equation (8) becomes

$$U_y(x, 0, t) = 2 \sigma_0 \int_0^t \sqrt{\frac{b}{a(\tau)-x}} K_n\left(\frac{a(\tau)}{b}\right) \frac{\partial a(\tau)}{\partial \tau} D_{cr}(t-\tau) d\tau \quad (9)$$

for $1 - \left| \frac{x}{a(t)} \right| < < 1$

When acceleration of the moving crack tip is small, one can assume

$$\frac{\partial a(t)}{\partial t} = V$$

where V is the crack velocity.

The variable change $\xi = Vt$ leads to the following form:

$$U_y(x, 0, v) = 2 \sigma_0 \int_x^{a_t} \sqrt{\frac{b}{\xi-x}} K_n(\xi/b) D_{cr}\left(\frac{a_t - \xi}{v}\right) d\xi \quad (10)$$

for $1 - \left| \frac{x}{a_t} \right| < < 1$

Substitute equation (2) and equation (10) into equation (1);

then one gets

$$-\Delta E_c = \int_{x=a}^{a+\Delta a} \left[\sigma_o \frac{K_n}{\sqrt{\frac{x-a}{b}}} \int_{x-\Delta a}^{a(t)} 2\sigma_o \sqrt{\frac{b}{\xi-(x-\Delta a)}} K_n\left(\frac{\xi}{b}\right) D_{cr}\left(\frac{a_t - \xi}{V}\right) d\xi \right] dx \quad (11)$$

By using spectrum form of the creep function

$$D_{cr}(t) = \frac{1}{E_r} \int_0^{\infty} L(\tau) \exp(-t/\tau) \frac{d\tau}{\tau}$$

and

$$\dot{E}_c = \frac{\Delta E_c}{\Delta a} V ,$$

equation (11) is reduced to the following form:

$$\dot{E}_c = 2\pi K_n^2 \sigma_o^2 b V D_{cr}\left(\frac{\Delta a}{V}\right) \quad (12)$$

where K_n is regarded as a constant for small intervals between a and $a+\Delta a$.

Taking the surface energy S and effect of temperature T ($^{\circ}K$) into consideration, equation (11) changes into

$$D_{cr}\left(\frac{\Delta a}{V a_T}\right) = \frac{1}{\pi} \frac{1}{b \sigma_o^2} \frac{S}{K_n^2(a/b)} \frac{T}{273}. \quad (13)$$

Equation (13) governs crack velocity for small Δa and small acceleration of crack velocity.

One can get the fracture criterion for viscoelastic materials by setting limit value of D_{cr} as $1/E_r$; E_r is the rubbery modulus.

$$\sigma_o \text{ crit.} = \sqrt{\frac{SE_r T}{\pi b K_n^2 (a/b) 273}} \quad (14)$$

2. EXPERIMENT

2.1 Apparatus and Specimen

A scheme of the method of this experiment is shown in Figure 4. A viscoelastic strip of width $2b$ is stretched between upper and lower edge with no initial strain as shown by the solid line. A small initial crack is located at the center. This initial crack opens and propagates to right and left direction by edge displacement which gives strain ϵ_o as shown by the dotted line.

Pictures of a propagating crack are taken by a motor-driven 35 mm camera which is controlled by a timing device.

Apparatus

Figure 5 shows a scheme of a mechanism which gives a certain strain to the specimen and an overall view of the apparatus. The lower edge is fixed to the base and the upper edge is supported by two screws. An induction motor drives the screws through a chain and gives a certain amount of edge displacement to the upper edge, edge displacement is regulated by a limit switch. The moving velocity of the upper edge is 0.05 in./sec., which is equivalent to the strain rate 3.74%/sec for the width of specimen; 1.375 inches.

Specimen

Chemical Property

Solithane 113

(Polyester Elastomer)

manufactured by the Thiokol Chemical Company. This material is a synthetic resin of two components, Urethane Resin and Catalyst. Mechanical

properties depend on the mixing ratio of the two components. Mixing ratio of 50%-50% in volume is employed in the experiment.

Mechanical Characteristics: Rubbery Young's modulus is about 440 psi at room temperature. Relaxation modulus is shown in Figure 6.

Viscoelastic properties do not have any influence on the loading process because the relaxation time is much smaller than the initial time [4] of crack propagation.

Geometric specification of the specimen is shown in Figure 7.

2.2 Condition of Experiment

Temperature of environment 23°C (73°F)

Experiment No.	Strain Applied	Interval Time of Each Picture
601	6.5%	10 sec.
701	7.5%	8 sec.
805-805	8.5%	2.5 sec.
901	9.5%	2.5 sec.

2.3 Results

The velocity of crack propagation is obtained by graphical differentiation of time-crack position plots. Figure 8 shows velocities versus crack tip position for each experiment. Six cases out of 8 experiments show maximum velocity between $a/b \approx 0.5$ and $a/b \approx 1.0$. The crack velocity reaches a certain value for $a/b > 1.5 \sim 2.0$. The crack velocity is very sensitive to change of strain.

3. COMPARISON BETWEEN THEORY AND EXPERIMENT

Theoretical Calculation

Equation (13) can be changed into

$$V = \frac{\Delta a}{a_T} / D_{cr}^{-1} \left[\frac{1}{\pi} \cdot \frac{1}{b \sigma_o^2} \cdot \frac{S}{K_n^2(a/b)} \cdot \frac{T}{273} \right] . \quad (15)$$

D_{cr}^{-1} denotes the inverse function of the creep function D_{cr} , that is

$$D_{cr}(A) = B$$

$$D_{cr}^{-1}(B) = A.$$

One can simplify equation (15) into a function of the stress intensity factor K_n and the strain ϵ_o when other values and creep function are specified.

$$V = \mathfrak{F}(\epsilon_o K_n) \quad (16)$$

Velocity is calculated and ascertained experimentally by H. K. Mueller [2] in the case of steady state crack propagation

in which the crack length is large enough for considering stress intensity factor K_n as a constant value.

Figure 9 shows the crack velocity at the steady state crack propagation for specified experiment in the previous section. This curve is calculated under assumption of

Characteristic length	$\Delta a = 10^{-8}$ inches
Surface energy	$S = 3.21 \times 10^{-2}$ lbs/in.
Rubbery relaxation modulus	$E_r \approx 443$ psi
Creep function	Figure 10

Now one can make an attempt to determine crack velocity by introducing equivalent strain ϵ_o^* .

$$\epsilon_o^* \approx \epsilon_o \frac{K_n}{0.346} \quad (17)$$

The number 0.346 is the value of the stress intensity factor K_n for $a/b > 1.5$ where K_n is considered as a constant value. K_n is shown in Figure 4. Velocity can be read from Figure 9 by using ϵ_o^* which is obtained through equation (17).

Experimental values have a certain scattering range around the theoretical curve, so it may be reasonable to set upper and lower limits such that most of the experimental values drop between these two limits in Figure 9.

The essential difficulty in this experiment is that the crack velocity is very sensitive against the condition of the experiment. The leading factors which cause experimental errors are:

- (1) Error of setting strain
- (2) Error of estimated temperature of specimen
- (3) Error of measured crack velocity through graphical differentiation.

One can reduce these errors by improvement method and apparatus of experiment.

It will be desirable to use statistical judgement for these experiments in which measured values scatter widely like this experiment.

COMPARISON BETWEEN THEORY AND EXPERIMENT

Figure 11 shows measured velocity and calculated velocity versus crack tip position for each strain. Measured velocities from this experiment lie within the upper and lower limit in the case of $\epsilon_0 = 6.5\%$ and 8.5% . The experiment of $\epsilon_0 = 7.5\%$ shows exceptionally high crack velocity. It is unclear to say whether velocity is reasonable or not in the case of $\epsilon_0 = 9.5\%$. These experimental values are plotted again in Figure 9. One can see that most of the measured values of this experiment correspond to that of steady state crack propagation and also one can say that unusual value in the case of 7.5% can be negligible because steady velocity at $a/b > 1.5$ is outside of the limit.

CONCLUSION

The equation (13) is essentially adequate to the steady state (constant crack velocity) crack propagation in a viscoelastic strip.

This experiment intends to escalate applicable area of the equation (13) such that one can use the same equation for a smaller range of crack length where stress intensity factor K_n is not a constant but a function of crack length/strip width ratio (a/b).

There are some assumptions for making use of the equation (13) for $a/b > 1.5$. One is that the acceleration of crack propagation is small and another is that the characteristic length Δa is small. This assumption makes an unsteady problem. This experiment suggests a significant possibility to make use of the equation (13) for smaller crack lengths.

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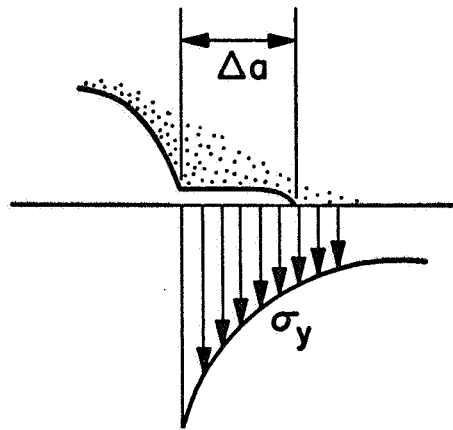
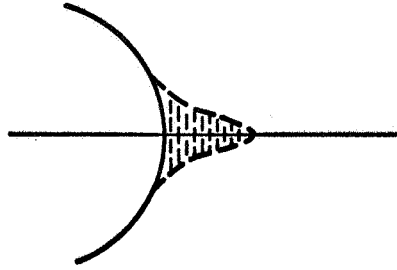


FIG. 1
MODEL OF CRACK TIP

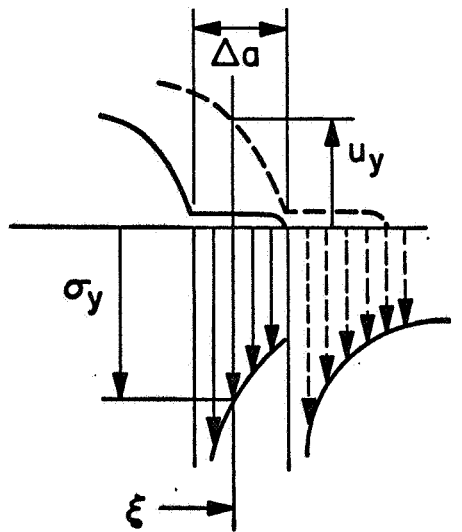


FIG. 2
OPENING CRACK

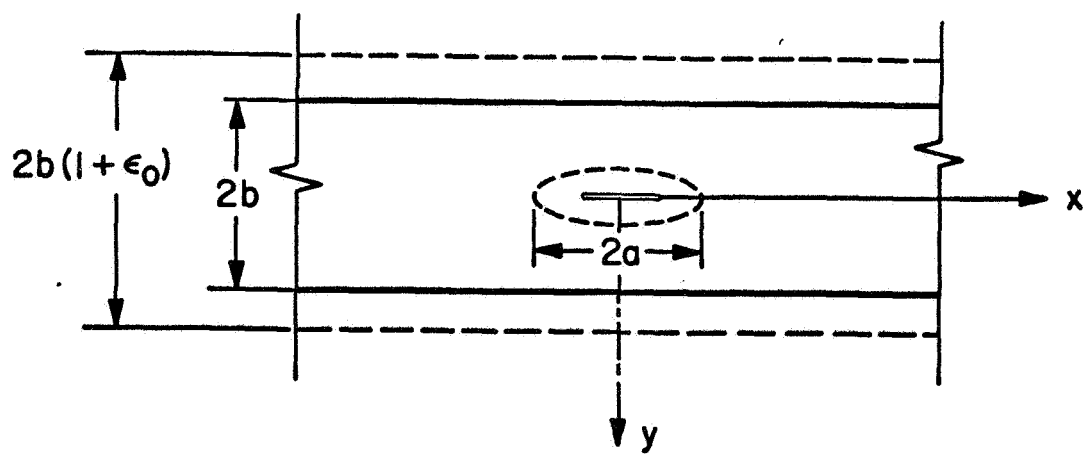


FIG. 3 STRIP GEOMETRY

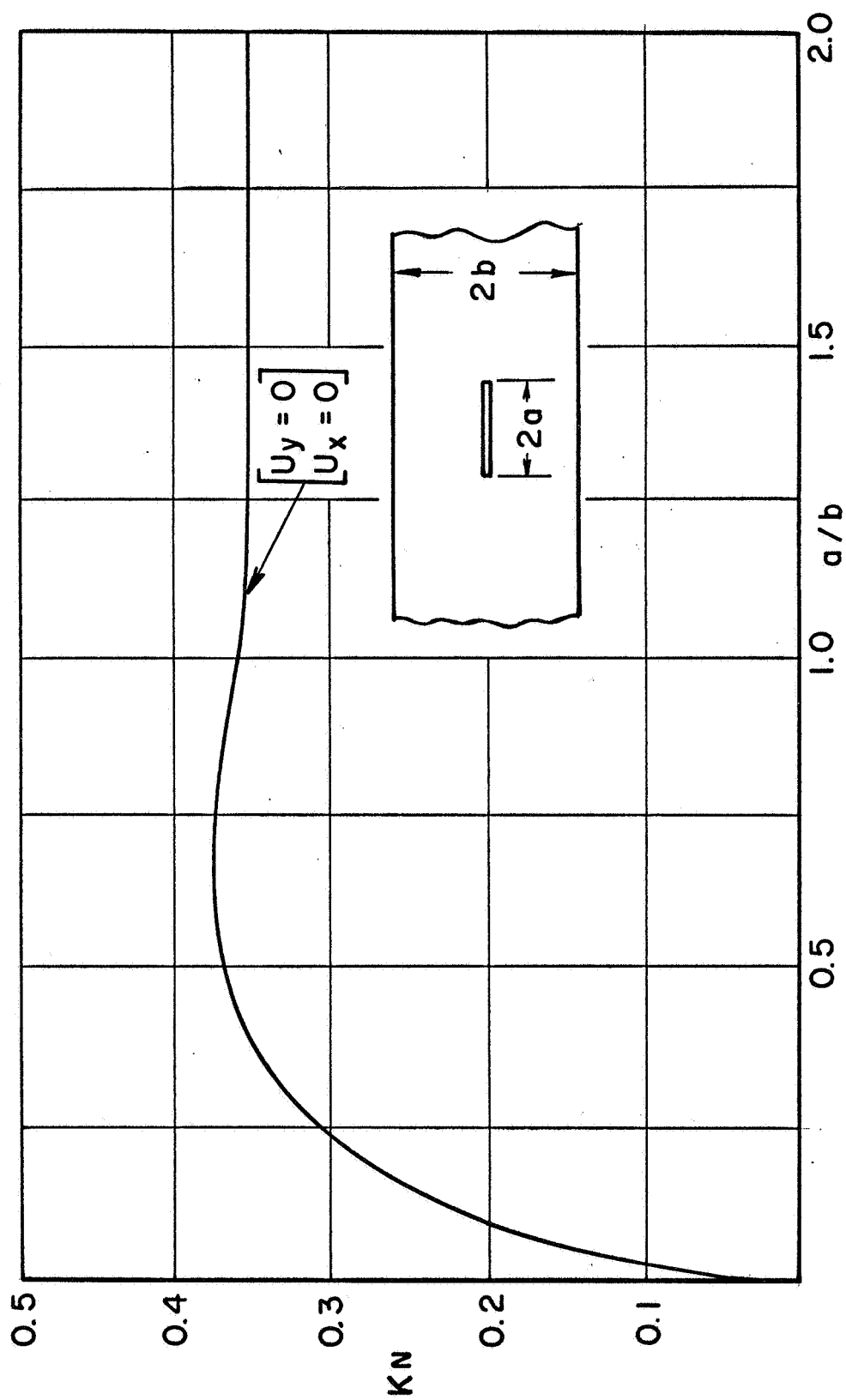


FIG. 4 STRESS INTENSITY FACTOR KN VERSUS CRACK TIP POSITION a/b

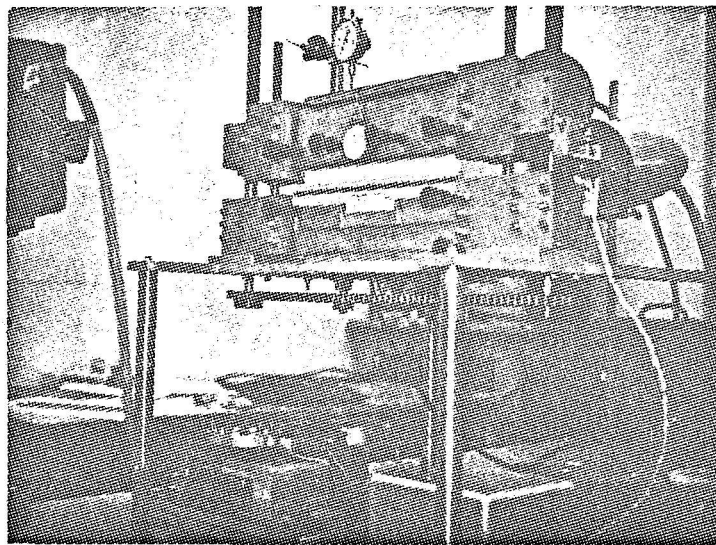
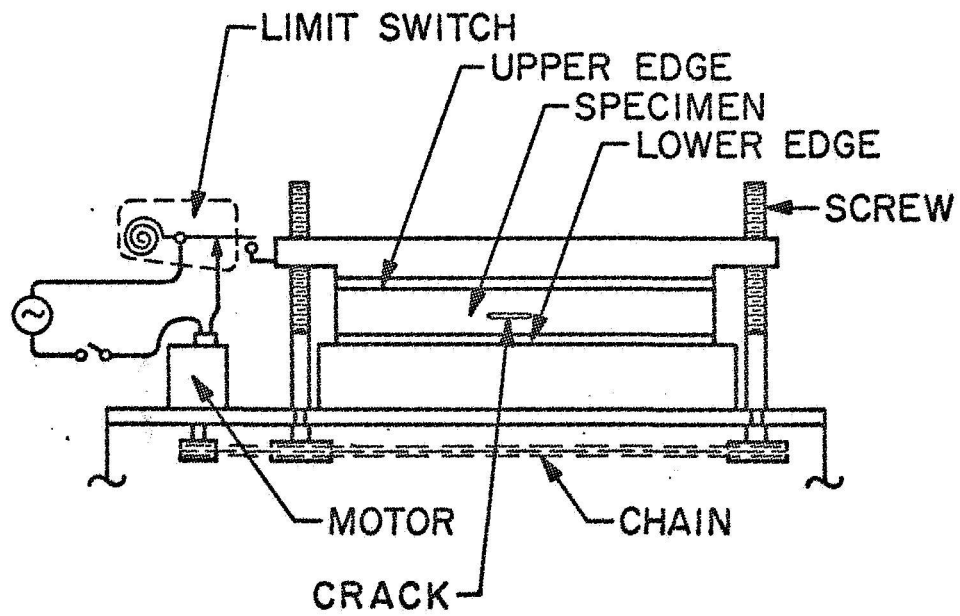


FIG. 5 APPARATUS

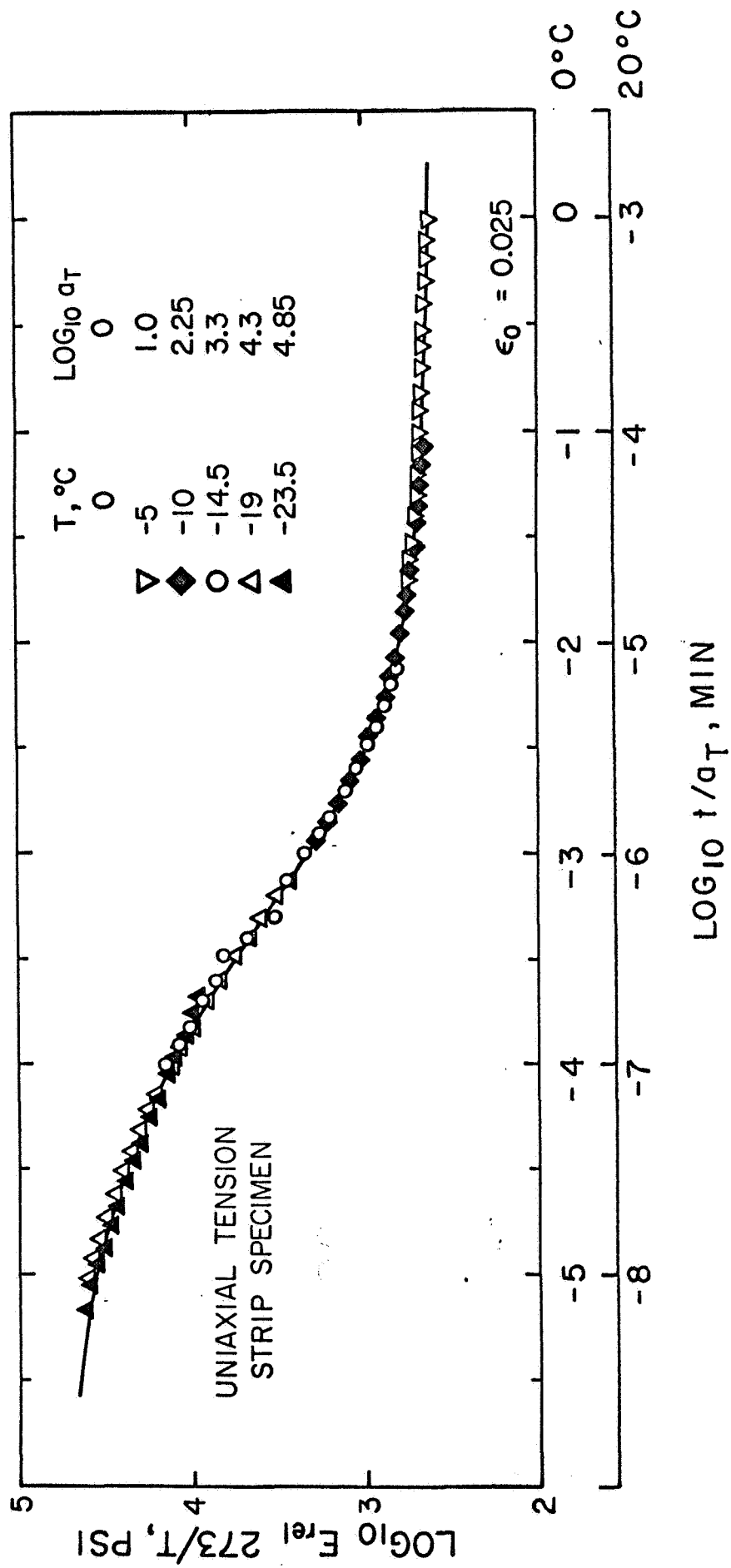


FIG. 6 RELAXATION MODULUS OF SOLITHANE 50/50

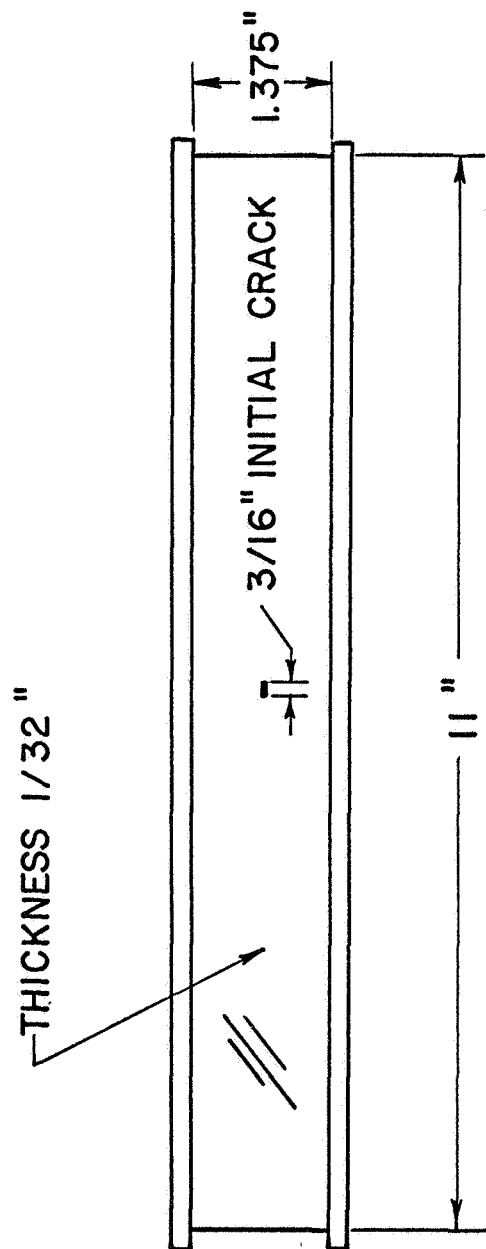


FIG. 7 SPECIMEN

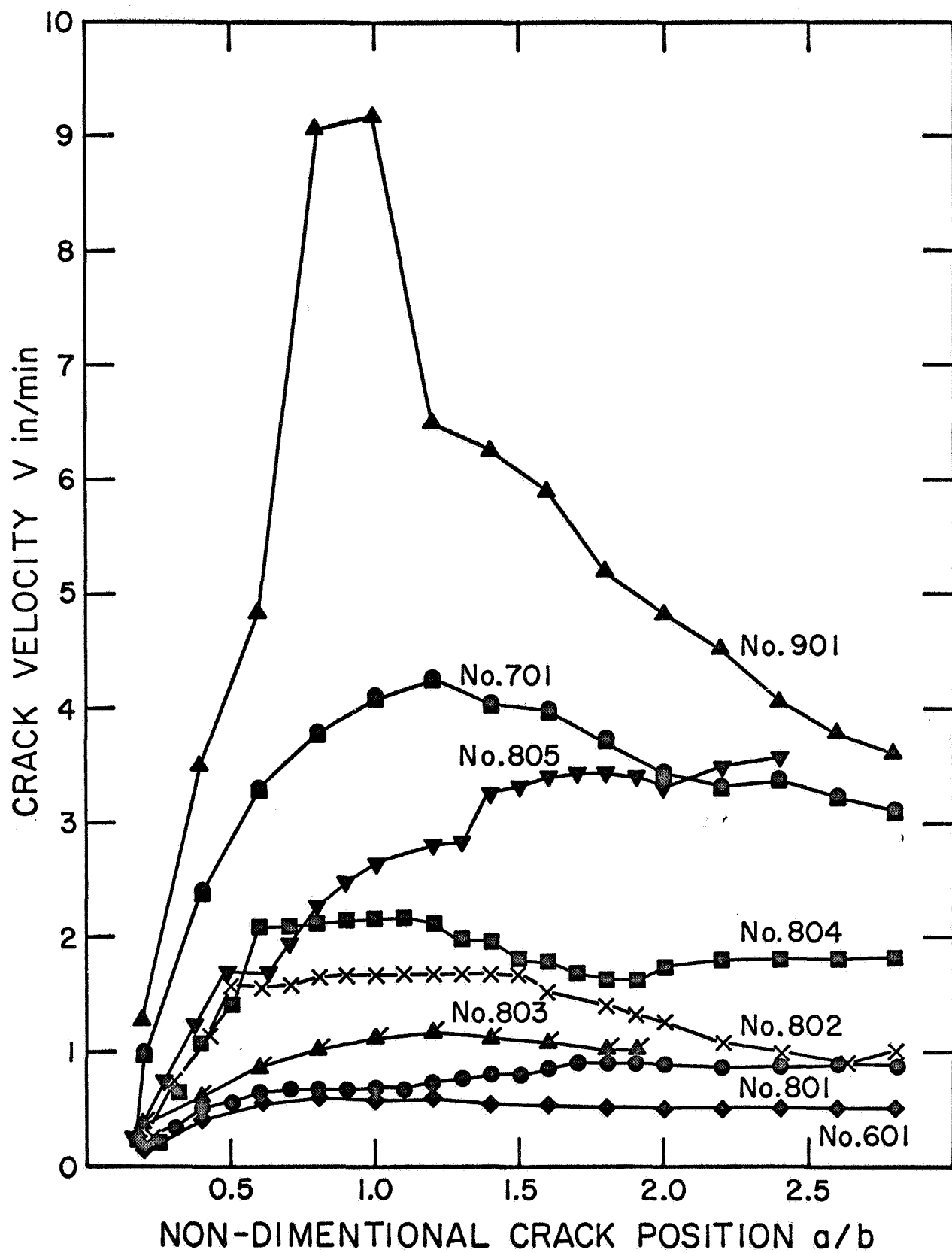


FIG. 8 RESULT OF EXPERIMENT

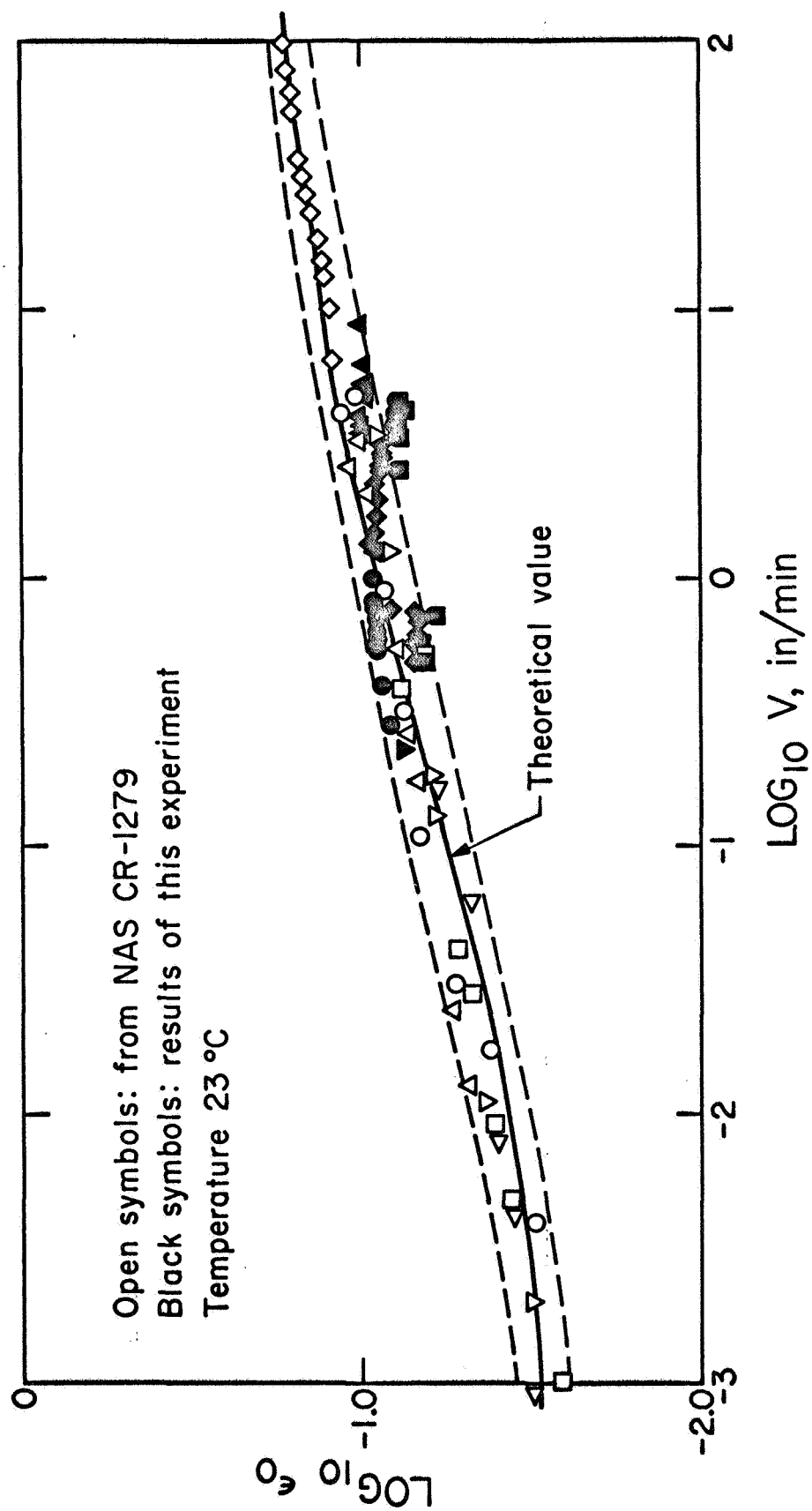


FIG. 9 CRACK VELOCITY VERSUS STRAIN IN STEADY STATE

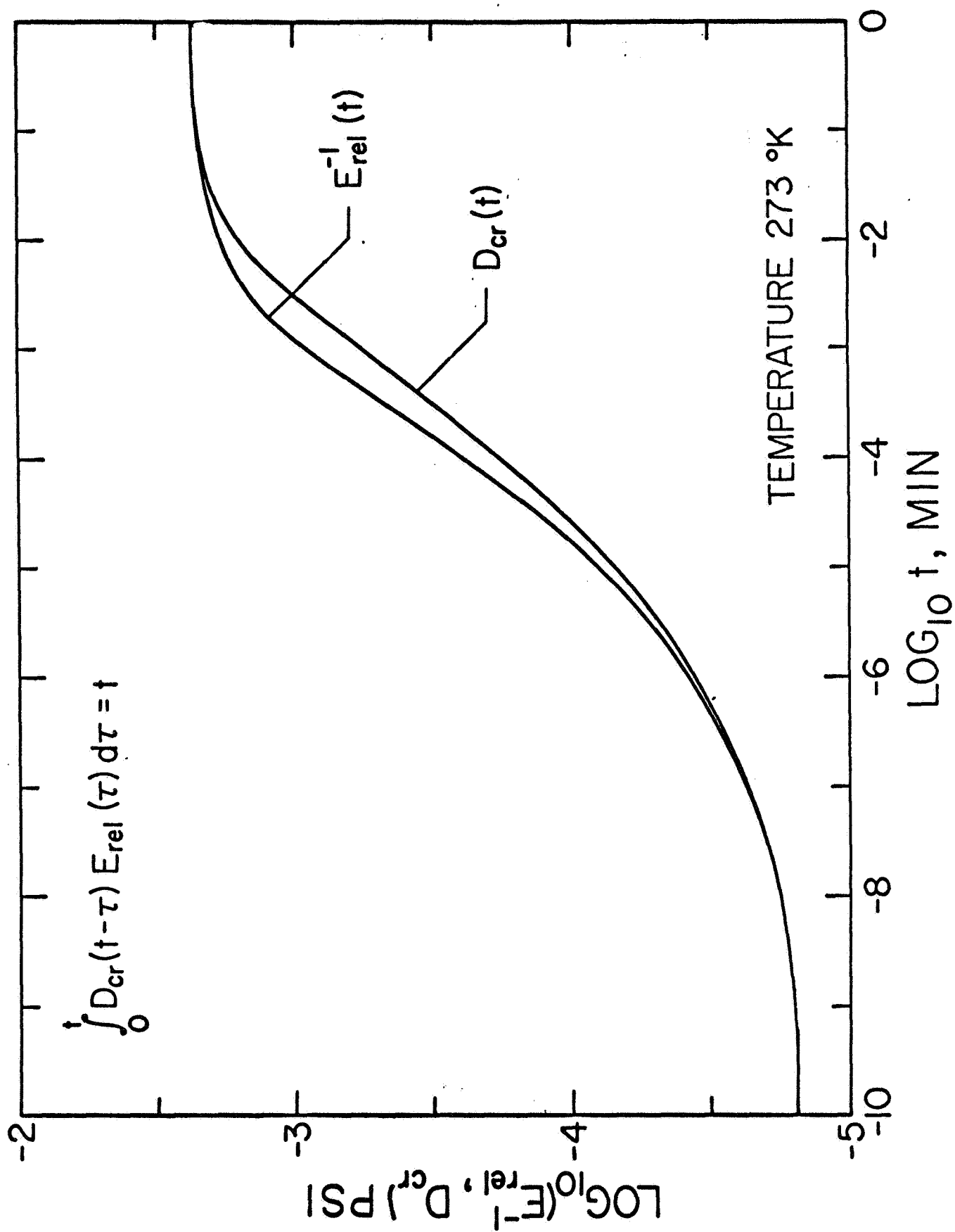


FIG. 10 THE RELAXATION FUNCTION, CREEP FUNCTION
 SOLITHANE 50/50

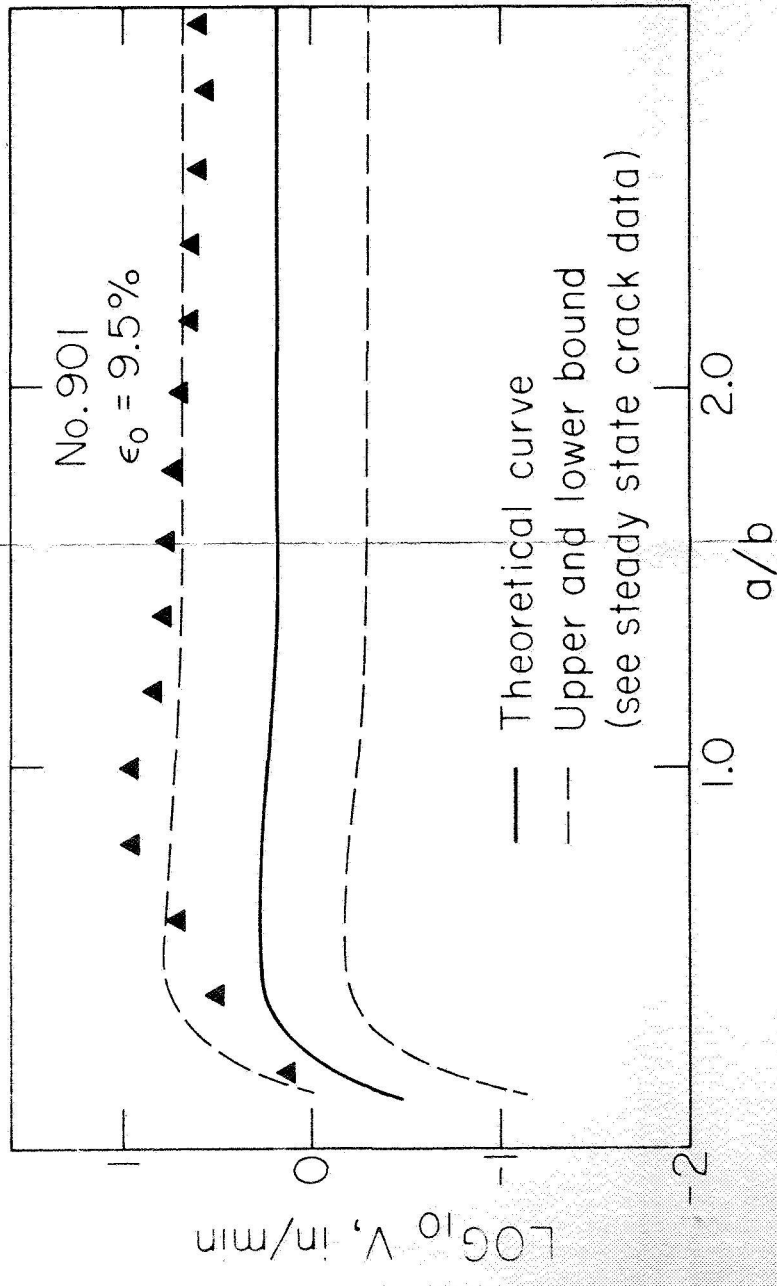
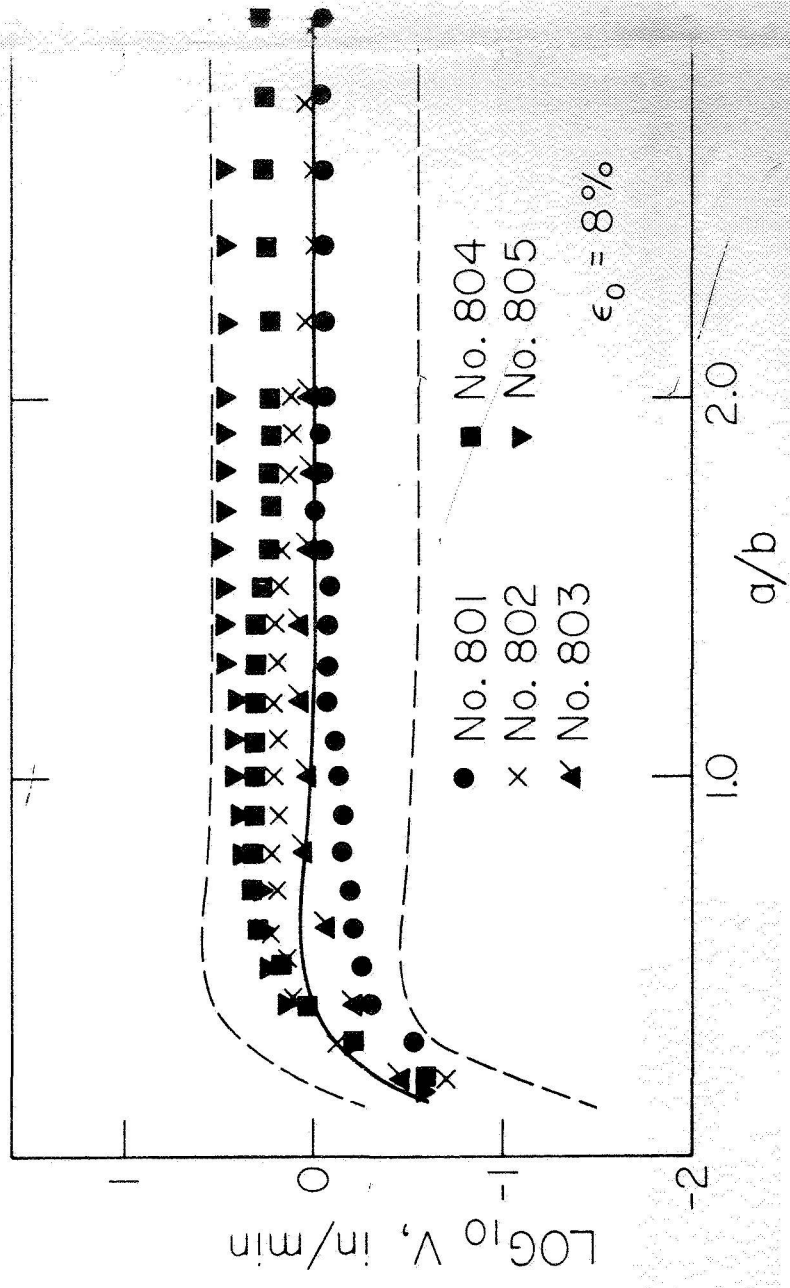
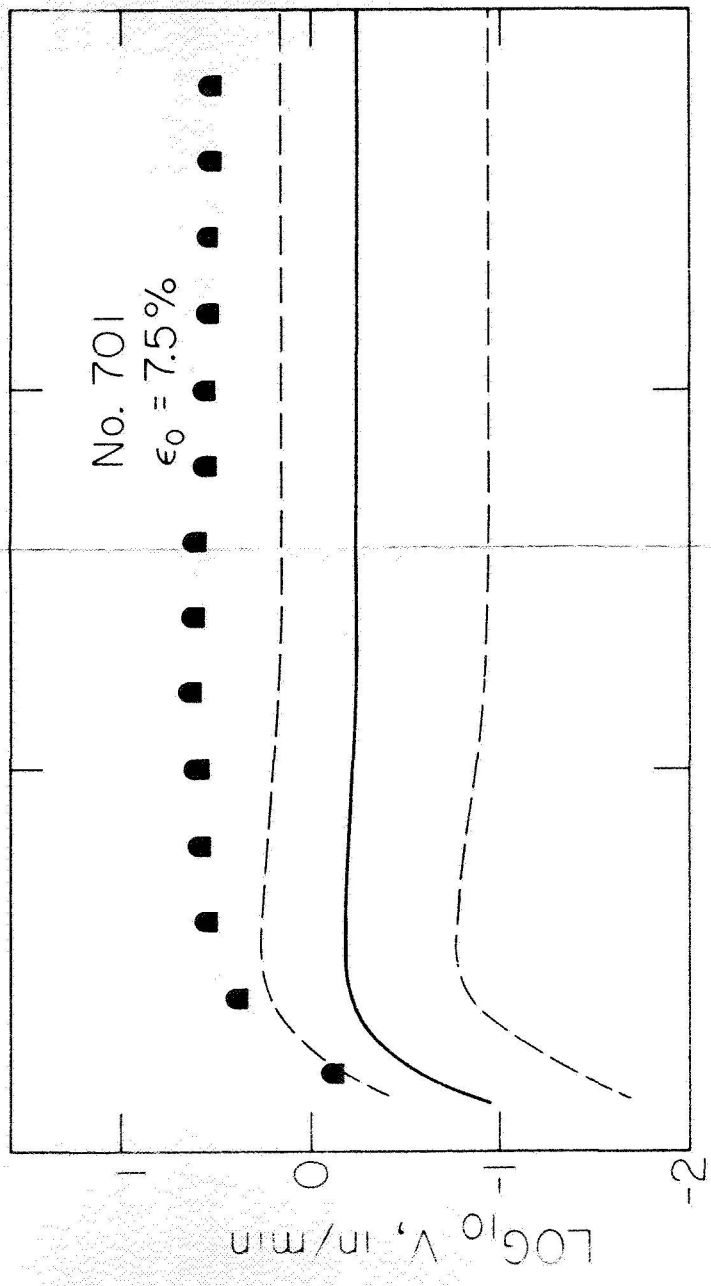
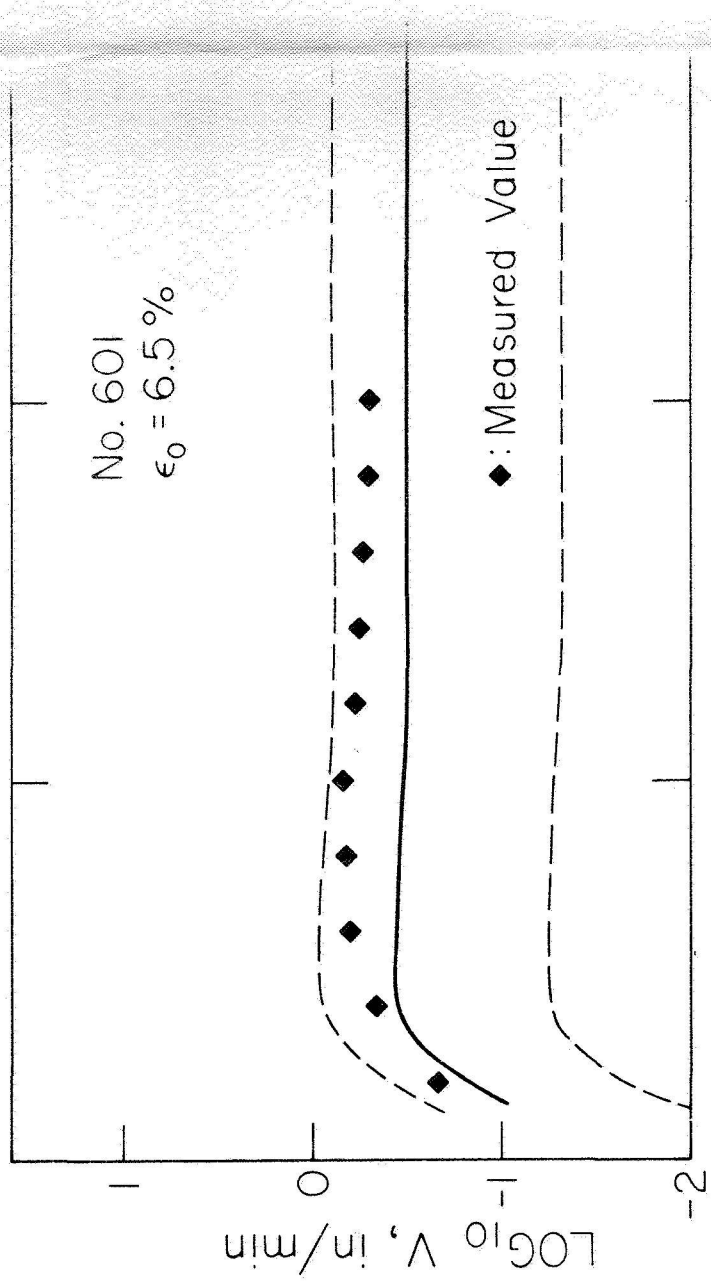


FIG. 11 MEASURED AND CALCULATED

CRACK TIP POSITION